



## Compressed sensing for resolution enhancement of hyperpolarized $^{13}\text{C}$ flyback 3D-MRSI

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### ABSTRACT

High polarization of nuclear spins in liquid state through dynamic nuclear polarization has enabled the direct monitoring of  $^{13}\text{C}$  metabolites in vivo at very high signal-to-noise, allowing for rapid assessment of tissue metabolism. The abundant SNR afforded by this hyperpolarization technique makes high-resolution  $^{13}\text{C}$  3D-MRSI feasible. However, the number of phase encodes that can be fit into the short acquisition time for hyperpolarized imaging limits spatial coverage and resolution. To take advantage of the high SNR available from hyperpolarization, we have applied compressed sensing to achieve a factor of 2 enhancement in spatial resolution without increasing acquisition time or decreasing coverage. In this paper, the design and testing of compressed sensing suited for a flyback  $^{13}\text{C}$  3D-MRSI sequence are presented. The key to this design was the undersampling of spectral k-space using a novel blipped scheme, thus taking advantage of the considerable sparsity in typical hyperpolarized  $^{13}\text{C}$  spectra. Phantom tests validated the accuracy of the compressed sensing approach and initial mouse experiments demonstrated in vivo feasibility.

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### 1. Introduction

Carbon-13 spectroscopy has traditionally been limited by low signal strength. With the development of techniques to maintain hyperpolarization of carbon-13 in liquid state [1], it has become possible to use  $^{13}\text{C}$  substrates (tracers) for medical imaging [2]. More recent studies have used the metabolically active substrate  $[1-^{13}\text{C}]$ pyruvate to examine its conversion to  $[1-^{13}\text{C}]$ lactate,  $[1-^{13}\text{C}]$ alanine, and  $^{13}\text{C}$ -bicarbonate [3–5]. Spectroscopic examination of these metabolic pathways in the presence and absence of disease has enormous diagnostic potential. Specifically, it has already been shown that the levels of  $^{13}\text{C}$  metabolic products differ between disease and non-disease states in a mouse model of prostate cancer [5]. As pointed out in [5], partial voluming may complicate the interpretation of non-disease spectra because of the small size of the normal mouse prostate. With the abundant SNR available in hyperpolarized studies, it would be beneficial to sacrifice some signal for improved spatial resolution, but the time limitation

imposed by T1 relaxation severely restricts the possible number of phase encode steps.

Recent advances in mathematical theory have opened the door for accurate reconstruction of sparse signals from sub-Nyquist sampling [6,7]. Less technical descriptions from the same authors, focusing on the practical limits of compressed sensing, have shown reconstructions from realistic data sets [8,9]. In addition, Lustig et al., in an exposition of the application of compressed sensing to MRI, shows that many MR images exhibit a high degree of sparsity and provides high quality proof of concept results drawn from multi-slice fast spin-echo brain imaging and 3DFT time of flight contrast-enhanced angiography [10]. Lustig lists three criteria for the successful application of compressed sensing: (1) the data have a sparse representation in a transform domain; (2) the aliasing from undersampling be incoherent in that transform domain; and (3) a non-linear reconstruction be used to enforce both sparsity of data and consistency with measurements. The successes in [10], coupled with the sparsity in hyperpolarized spectra, make hyperpolarized  $^{13}\text{C}$  spectroscopic imaging a logical choice for the application of compressed sensing. In other words, for hyperpolarized  $^{13}\text{C}$ , the first criterion is satisfied by the inherent sparseness of hyperpolarized spectra, and the third criterion can be met by using the same non-linear reconstruction from [10]. The primary challenge

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then is to satisfy the second criterion: developing an undersampling approach to achieve suitable incoherent aliasing.

The sparsity of hyperpolarized spectra has been previously exploited for accelerated imaging [11,12] by using *a priori* knowledge of metabolite resonance locations and linewidths (factor of 4 acceleration reported in a phantom demonstration [11]). Compressed sensing, however, requires no assumptions except that the underlying data are sparse in some domain.

A 3D-MRSI sequence using flyback echo-planar readout gradients [13,14] provides a fast method for acquiring hyperpolarized spectra and is less sensitive to timing errors, eddy currents, and  $B_0$  inhomogeneity than EPI and spiral readout schemes, especially in vivo. This paper presents a methodology for accelerating the acquisition of hyperpolarized spectra by a factor of 2 using compressed sensing in conjunction with a modified flyback echo-planar 3D-MRSI sequence. A few initial phantom and in vivo examples are presented as proof of concept.

## 2. Theory

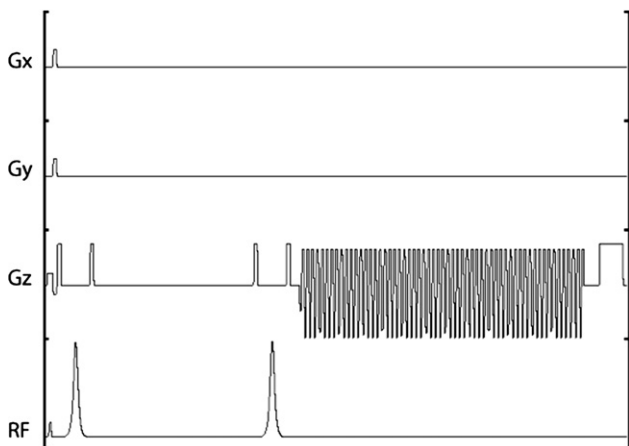
### 2.1. Pulse sequence

The pulse sequence developed for this study builds on the one diagrammed in Fig. 1, which is a double spin-echo sequence with a flyback echo-planar readout [14]. Details of the flyback design are given in [13]. As explained in [14], a spin-echo sequence was desired in order to mitigate the effects of  $B_0$  inhomogeneity and allow for a full echo acquisition. Owing to its insensitivity to transmit gain, this double adiabatic sequence outperforms a conventional spin-echo sequence in preserving hyperpolarization over the repeated excitations needed for the phase encodes in an MRSI acquisition [14]. Ultimately, this sequence reads out a rectilinear k-space trajectory, with a typical result being  $59 \times 8 \times 8 \times 16$  ( $k_f$ - $k_x$ - $k_y$ - $k_z$ ) 4D k-space data.

### 2.2. Key results from compressed sensing literature

Fundamentally, compressed sensing claims to perfectly reconstruct sparse signals of length  $N$  from a subset of samples. For example, suppose a length  $N$  discrete signal  $f$  consists of  $M$  non-zero points. Then, with extremely high probability,  $f$  can be recovered exactly from  $K$  Fourier measurements where

$$K \geq \text{Const} \cdot M \log N \quad (1)$$



**Fig. 1.** Double spin-echo sequence timing diagram. The RF consists of a small tip excitation followed by two adiabatic pulses (phase channel not shown). Phase encoding is along  $x$  and  $y$  while the 59-lobe flyback readout is along  $z$ . An echo is formed during the middle of the flyback readout with TE = 140 ms.

and the solution is found by solving the convex minimization problem

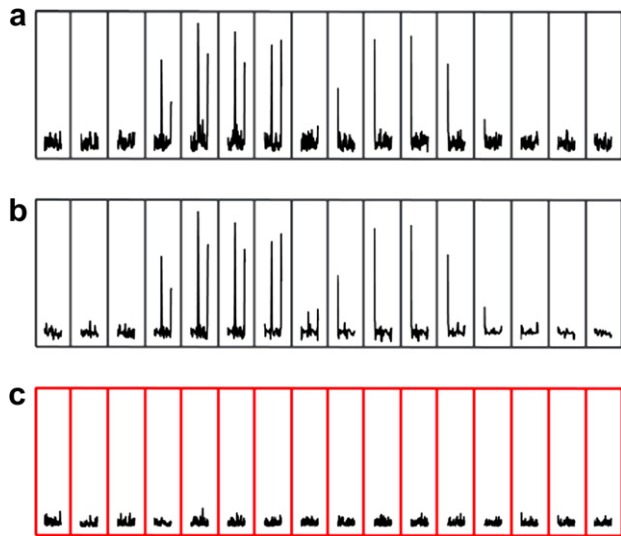
$$\min \sum_n |g[n]| \quad \text{s.t. } F_k \{g[n]\} = y \quad (2)$$

where  $F_k$  is the Fourier transform evaluated at  $K$  locations and  $y$  is the set of  $K$  measured Fourier coefficients [8]. In words, Eq. (2) states that for all reconstructions  $g[n]$  whose Fourier coefficients match those at the  $K$  measured positions, the unique and correct solution is the one that minimizes the absolute sum of  $g$ , i.e. the  $\ell_1$  norm in the object domain. The theorems of compressed sensing are actually much more general than this concrete example suggests. In other words, the signal  $f$  only needs to be sparse in some domain, not necessarily the object domain, and the  $K$  measurements do not necessarily have to be Fourier measurements. From a practical standpoint, the application specific values for  $M$ ,  $N$ , and the constant multiplier determine the feasibility of compressed sensing. Additionally, a real-world signal will never consist of just  $M$  non-zero points in any domain, but it will usually be well approximated by  $M$  sparse transform coefficients. For example, the fidelity with which compressed sensing reproduces an  $M$ -term wavelet approximation, i.e. the sparse domain being the wavelet domain, could serve as a benchmark for real-world signals such as NMR spectra [9]. For various  $N = 1024$  test signals in [8], Candes empirically determined that for compressed sensing to match the accuracy of an  $M$ -term wavelet representation,  $K \approx 3M - 5M$  measurements were required, which was also observed in [10]. As mentioned previously, for the actual implementation of compressed sensing, the  $K$  measurements must be collected with a sampling pattern that produces incoherent aliasing in the domain, such as the wavelet domain, where the signal shows sparsity [10]. A random sampling pattern in k-space almost always meets this criterion.

At this point, it is interesting to consider the connection between compressed sensing and existing techniques in NMR, such as maximum entropy [15,16] and minimum area [17] reconstruction, used for the related problem of computation of spectra from short, noisy data records. Recently, Stern et al. showed that a specific form of iterative thresholding, a technique similar to maximum entropy and minimum area reconstruction, is equivalent to the minimum  $\ell_1$  norm reconstruction in compressed sensing [18]. Additionally, Stern explains how  $\ell_1$  norm reconstruction gives insight into the performance of maximum entropy and minimum area reconstruction. Thus, compressed sensing could be viewed as a generalization of existing NMR techniques.

### 2.3. 3D-MRSI Signal

The most straightforward application of compressed sensing to hyperpolarized 3D-MRSI would be to undersample in  $k_x$  and  $k_y$ . For example, to achieve  $16 \times 8$  spatial resolution in the time of  $8 \times 8$  phase encodes, i.e. a speedup factor of 2, one could simply collect  $8 \times 8$  of the phase encodes in a conventional  $16 \times 8$  scan ( $K = 64$ ,  $N = 128$ ). However, our wavelet simulations have shown that such a small  $N$  leads to a relatively large  $M$  and thus does not provide enough sparsity to exploit. A better strategy would be to attempt undersampling in the  $k_f$  and  $k_x$  dimensions, considering that typical hyperpolarized acquisitions are inherently sparse in the spectral dimension. As shown in the wavelet simulations of Fig. 2, a signal of this type (spectral dimension and one spatial dimension) exhibits considerable sparsity. (Note that wavelet transforms were chosen because they do a good job of sparsifying NMR spectra [9], though other choices are possible as well.) The key point is that the majority of the sparsity occurs in the spectra and therefore the time domain

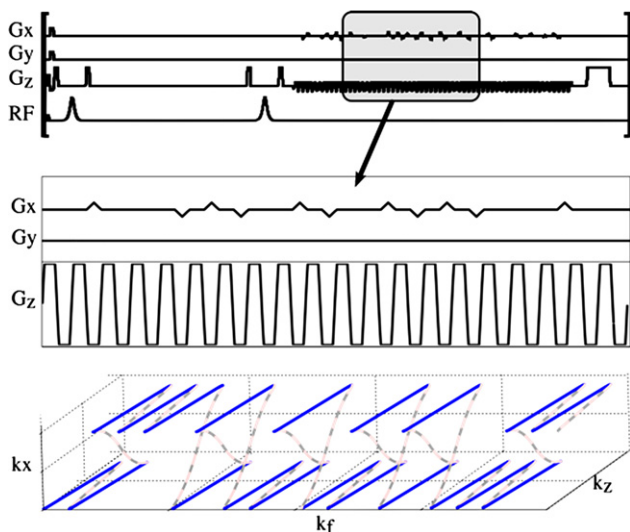


**Fig. 2.** Demonstration of wavelet compressibility of a  $^{13}\text{C}$  spectroscopic signal. A row of magnitude spectra ( $64 \times 16$ ) from a 3D-MRSI phantom data set (see Fig. 5 for examples of rows of spectra) was taken as the test signal. Note that the 59 spectral points from the 59 flyback lobes were zero-padded to 64 because the wavelet software we used required dyadic numbers. (a) The 16 original spectra. (b) A 2D Daubechies wavelet transform was applied to the  $64 \times 16$  data, after which the top 10% wavelet coefficients were retained and the inverse 2D wavelet transform taken. (c) The magnitude error between (a) and (b). Note that (a)–(c) have the same y-axis scale. The  $64 \times 16$  data were reconstructed very accurately from only 10% of their wavelet coefficients, showing that the signal of interest exhibits considerable fundamental sparsity.

should be undersampled. However, the implementation of time domain undersampling is not at all straightforward, as the next section demonstrates. A scheme to undersample in the time domain as well as in one spatial domain was employed, mainly exploiting spectral sparsity but some spatial sparsity as well.

#### 2.4. Implementation of $k_f$ - $k_x$ incoherent sampling

The key to implementing a k-space trajectory that randomly undersamples in  $k_f$ - $k_x$  lies in the random sampling of  $k_f = t$  using blips. Figs. 3 and 4 illustrate a scheme that achieves  $k_f$  sub-sam-



**Fig. 3.** Blipped scheme for  $k_f$ - $k_x$  sub-sampling. (Top) The only modification to the pulse sequence shown in Fig. 1 is the addition of blips during the rewind portions of the flyback readout. The area of each blip is the area in an  $x$ -phase encode step. Bottom: Associated order of k-space readout. A single readout now covers two  $k_f$ - $k_x$  lines.

pling by hopping back and forth between adjacent  $k_x$  lines during a flyback readout. In this manner, data from two  $k_f$ - $k_x$  lines are acquired during a single phase encode, in effect randomly undersampling in time. Thus,  $16 \times 8$  resolution can be achieved in half the time by collecting  $8 \times 8$  of the readouts in a conventional  $16 \times 8$  scan. This approach is somewhat similar to the  $k$ - $t$  sparse scheme in [19,20], but here we apply gradients to move around in  $k_f$  space instead of reordering phase encodes. This blipped scheme addresses the design challenge of generating sufficient incoherent aliasing through random undersampling, in other words meeting the second criterion for the successful application of compressed sensing. Without the blips, there would be too much structure to the undersampling, which would lead to coherent aliasing. To reiterate, the design in Fig. 3 achieves twofold undersampling by jumping between two lines. A design to achieve threefold undersampling would have to jump between three lines, and a design to achieve fourfold undersampling would have to jump between four lines. Finally, we used the Duyn method [21,22] to measure the actual k-space trajectory traced out by our blips. As expected, on a modern clinical scanner with eddy current compensation, the measured k-space trajectory closely matched the intended one. In other words, the blips produced minimal side effects and unintended k-space deviations were negligible.

### 3. Experimental

#### 3.1. Pulse sequence and hardware

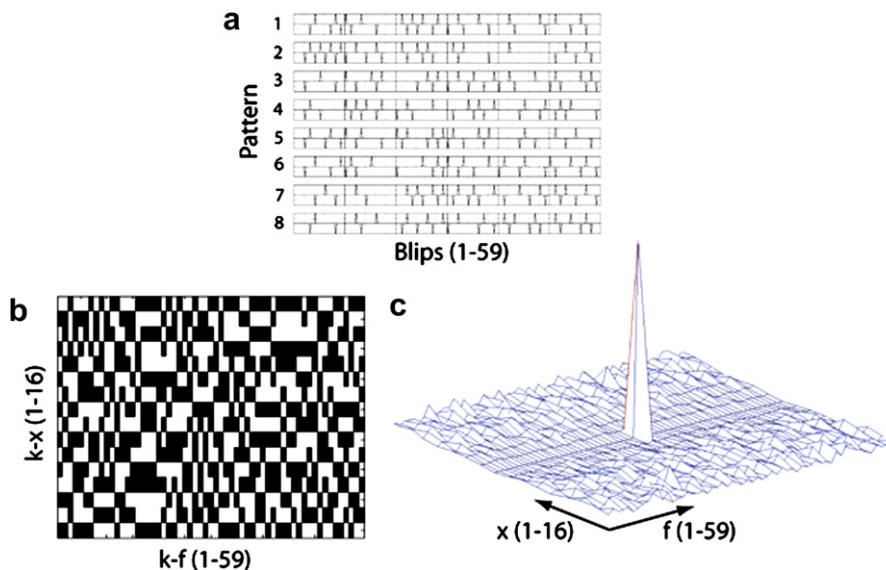
The source code from [14], originally a free induction decay (FID) MRSI sequence, was modified to incorporate triangular gradient blips. In an attempt to minimize eddy current effects, the blips were made 0.8 ms, relatively wide considering the time between adjacent flat flyback portions was 1.16 ms. The amplitude of the blips was calculated by the source code so that each blip's area equaled the area in a phase encode increment. As in [5], a variable flip angle (VFA) scheme [23], i.e. increasing flip angle over time to compensate for the loss in hyperpolarized signal, was used in the in vivo experiments. The actual  $n$ th flip angle  $\theta[n]$  precalculated by the source code for a given acquisition of  $N$  flips was as follows:

$$\theta[n] = \begin{cases} 90^\circ & \text{if } n = N \\ \arctan(\sin(\theta[n+1])) & \text{if } n < N \end{cases} \quad (3)$$

For example, in our acquisition with  $N = 8 \times 8 = 64$  readouts,  $\theta[64] = 90^\circ$ ,  $\theta[63] = \arctan(\sin(90^\circ)) = 45^\circ$ ,  $\theta[62] = \arctan(\sin(45^\circ)) = 35.3^\circ, \dots, \theta[1] = 7.2^\circ$ . Calibration of the pulse angles was performed on the day of each study using a prescan of a corn oil phantom. In addition, for the in vivo experiments, as in [5], reordering of phase encodes to collect data near the k-space origin first was also employed. For all experiments, T2-weighted images were acquired with a fast spin-echo sequence, after which MRSI data, phase encode localized in  $x/y$  with flyback readout in the S/I direction  $z$ , were collected. All experiments were performed on a General Electric EXCITE 3T (Waukesha, WI) clinical scanner equipped with 40 mT/m, 150 mT/m/ms gradients and a broadband RF amplifier. Custom-built, dual-tuned  $^1\text{H}/^{13}\text{C}$  transmit/receive coils were used for all phantom and animal experiments.

#### 3.2. Reconstruction

For acquisitions without blipped gradients, the reconstruction procedure, carried out with custom MRSI software [24], was as follows: (1) sample the raw flyback data to obtain a 4D matrix of k-space data; (2) apodize each FID and apply a linear phase correction to the spectral samples as described in [13] to account for the tilted k-space trajectory characteristic of a flyback readout; and (3)



**Fig. 4.** Blipped patterns to cover 16  $k_f$ - $k_x$  lines, resulting coverage, and point spread function. (a) Actual eight blipped patterns used to cover 16  $k_f$ - $k_x$  lines in a pseudo-random manner. (b) Associated  $k$ -space sampling. Because twice as much  $k$ -space is covered in the time of eight phase encodes, half of the  $59 \times 16$   $k_f$ - $k_x$  points are missing (missing points are black). (c) 2D point spread function of pseudo-random pattern in (b).

perform a 4D-Fourier transform with zero-padding of the spectra. For blipped acquisitions, the processing pipeline was modified in that the flyback sampling was performed in MATLAB (Mathworks Inc., Natick, MA) and the  $k$ -space points missed by the blipped trajectory were iteratively filled in using a non-linear conjugate gradient implementation of Eq. (2) [10]. Specifically, the reconstruction procedure for blipped acquisitions was as follows: (1) order the raw blipped flyback data to obtain a 4D matrix of  $k$ -space data missing half of its  $k_f$ - $k_x$  points; (2) inverse Fourier transform the fully sampled  $k_y$  and  $k_z$  dimensions; (3) iteratively fill in the missing  $k_f$ - $k_x$  points in the 4D matrix using the algorithm from [10]; (4) forward Fourier transform the  $k_y$  and  $k_z$  dimensions to put the data back into the form of a filled  $k_f$ - $k_x$ - $k_y$ - $k_z$  set; (5) apodize each FID and apply a linear phase correction; and (6) perform a 4D-Fourier transform with zero-padding of the spectra. Transforming in the fully sampled  $k_y$  and  $k_z$  dimensions allowed us to separate the multi-dimensional reconstruction problem into many separable 2D reconstructions, reducing the memory requirements and allowing parallel processing as was done for the 3D angiography example in [10]. The total reconstruction time for the normal and compressed sensing reconstructions were  $\sim 5$  s and  $\sim 20$  min, respectively, on a 1-GHz, 2 GB RAM Sun workstation running Red Hat Linux. The compressed sensing reconstruction was implemented in MATLAB. We expect significant speed improvement with code optimization.

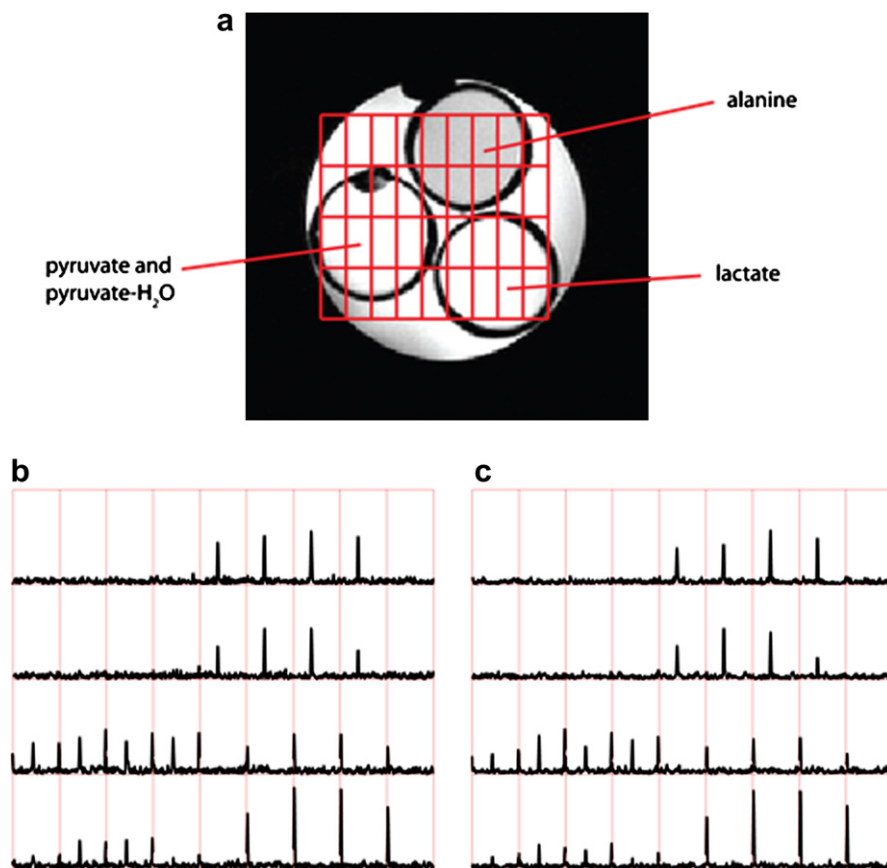
### 3.3. Phantom

Experiments on a cylindrical phantom (Fig. 5a) ( $n=3$ , with repositioning for separate trials) containing  $^{13}\text{C}$ -labeled pyruvate/pyruvate- $\text{H}_2\text{O}$ , lactate, and alanine in three respective inner spheres, were performed to verify the accuracy of the compressed sensing reconstruction. For both unblipped and blipped acquisitions, a flip angle of  $10^\circ$ , TE = 140 ms, TR = 2 s, FOV = 8 cm  $\times$  8 cm, and  $16 \times 8$  resolution were used. The  $16 \times 8$  unblipped acquisition with the standard reconstruction served as the gold standard. For the  $16 \times 8$  blipped acquisition, acquired in half the time, the modified processing pipeline as discussed in the previous section was used. The sparsifying transform was a 1D length-4 Daubechies wavelet transform in the spectral dimension, meaning the algo-

rithm presumed sparsity of the spectral peaks and tried to minimize the  $\ell_1$  norm of a wavelet transform of the  $k_f$  data. In addition, as is commonly done [10], a total variation (TV) penalty was added to promote sparsity of finite differences. The weights given to the wavelet transform and TV penalty, and thus the amount of denoising and data fidelity, were selected manually by testing a few values on one phantom acquisition. The same weights were used for subsequent phantom and animal experiments. Specifically, using the software described in [10], which normalizes the maximum signal in the object domain to 1, the TV penalty and transform weights were both 0.01.

### 3.4. Mouse

We performed normal and compressed sensing comparisons for three separate mice. For the *in vivo* experiment whose results are shown in Fig. 6, a prototype DNP polarizer developed and constructed by GE Healthcare (Malmö, Sweden) was used to achieve  $\sim 23\%$  liquid state polarization of [ $^{13}\text{C}$ ]pyruvate. Due to unavailability of the prototype machine for the second and third *in vivo* comparisons, a Hypersense<sup>TM</sup> DNP polarizer (Oxford Instruments, Abingdon, UK), which is a commercial version of the prototype machine, was used to achieve polarizations of  $\sim 11\%$  and  $\sim 18\%$ , respectively. The polarization was measured by extracting a small aliquot of the dissolved solution and measuring its FID intensity with a custom low-field spectrometer.  $\sim 300$   $\mu\text{L}$  ( $\sim 80$  mM) samples were injected into a surgically placed jugular vein catheter of a Transgenic Adenocarcinoma of Mouse Prostate (TRAMP) mouse within  $\sim 20$  s of dissolution. The particular TRAMP mouse for the first trial had a large prostate tumor with many relatively homogeneous tumor voxels across the FOV, making quantitative comparisons easier. The other two mice had smaller, yet still relatively homogeneous, tumors. For each trial, two runs were done ( $\sim 2$  h apart), once for an unblipped  $59 \times 8 \times 8 \times 16$  standard acquisition and again for a blipped  $59 \times 16 \times 8 \times 16$  compressed sensing acquisition. The acquisition parameters for both runs were TE = 140 ms, TR = 215 ms (total acquisition time of 14 s), variable flip angle activated, reordered phase encodes, and FOV = 4 cm  $\times$  4 cm. The blipped acquisition, using  $8 \times 8$  of the readouts from a conventional  $16 \times 8$  scan, was acquired after the



**Fig. 5.**  $16 \times 8$  phantom comparison of normal vs. undersampled. (a) T2-weighted image of  $^{13}\text{C}$  phantom done before spectral acquisitions. (b) Spectra from normal, unblipped acquisition corresponding to the highlighted voxels from (a). (c) Spectra from compressed sensing reconstructed, blipped acquisition corresponding to the highlighted voxels from (a).

unblipped one. All animal studies were carried out under a protocol approved by the Institutional Animal Care and Use Committee. A more detailed description of the polarization and animal care procedures can be found in [4,5].

## 4. Results and discussion

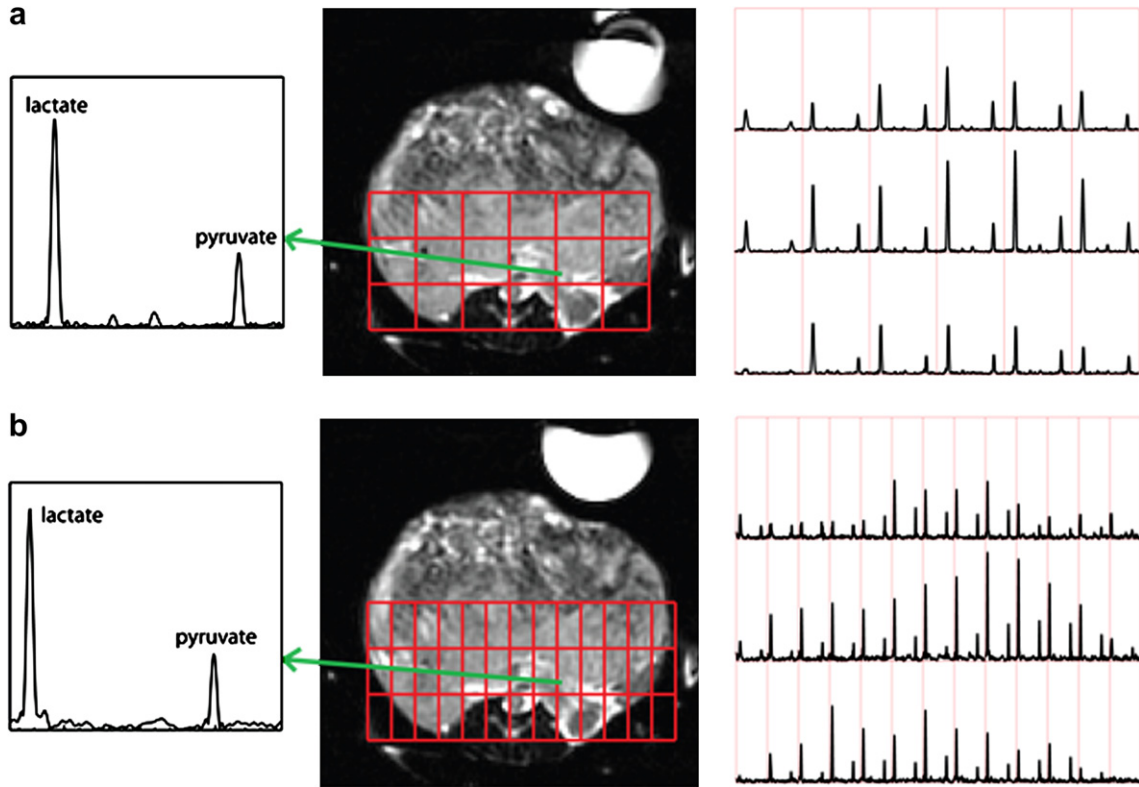
### 4.1. Phantom

Fig. 5 shows a side-by-side comparison of a slice of final processed spectra from representative  $16 \times 8$  unblipped and blipped acquisitions (corresponding to Trial 1 in Table 1), with the blipped acquisition taken immediately after the unblipped one. Qualitatively, the spectra match up extremely well. Table 1 gives a quantitative comparison of the two reconstructions for the three trials, listing SNR and metabolite peak ratios. The accelerated acquisition ratios were always within 10% of those from the fully sampled acquisitions, which was about the same as the accuracy reported in [11]. Because we typically draw biological conclusions from final processed spectra, SNR was calculated with the magnitude spectra after apodization and zero-padding. Typically, halving the scan time, as was done for the blipped acquisition, would reduce SNR by a factor of square root of 2. Due to the denoising properties of the compressed sensing reconstruction combined with apodization, the SNR did not drop. The  $\ell_1$  penalty used in compressed sensing is essentially a denoising procedure, also referred to in the literature as basis-pursuit denoising [25] and is closely related to wavelet denoising

schemes [26–28]. The  $\ell_1$  reconstruction transforms the signal into a domain in which the signal exists in only a few significant coefficients, whereas noise resides in the majority of coefficients, and filters the noise by heavily penalizing the small coefficients. In addition, in compressed sensing, the undersampling itself generates incoherent aliasing which appears as noise [10], which is penalized and filtered. Thus, the process of the  $\ell_1$  reconstruction picking a solution to the compressed sensing problem, in other words the underdetermined problem caused by undersampling, has the natural side effect of denoising. Therefore, the compressed sensing SNR is controllable in the sense that merely adjusting the  $\ell_1$  denoising parameters in the reconstruction would lead to higher SNR. However, too much denoising could lead to metabolite peak height distortion manifested by underestimating the true peak height. In this study, we tested denoising parameters within an order of magnitude as those in [10] and selected ones that performed well. As for the peak ratio calculations, the metabolite peak heights used were average peak heights over voxels with little or no partial voluming. The metabolite ratios for normal and compressed sensing data sets were similar, suggesting compressed sensing could be compatible with metabolite quantitation.

### 4.2. Mouse

Fig. 6 shows a comparison of  $8 \times 8$  mouse tumor data from a conventional scan and data from the same mouse acquired  $\sim 2$  h later with a  $16 \times 8$  compressed sensing acquisition. Qualitatively, the two sets of data appear similar, both showing elevated lactate



**Fig. 6.** Comparison of  $8 \times 8$  normal mouse data and  $16 \times 8$  undersampled mouse data in a region of prostate tumor. (a) Normal  $8 \times 8$  data. The left shows the spectrum with the highest lactate peak, the middle shows the T2-weighted anatomical image, and the right shows spectra highlighted in the anatomical image. (b) Corresponding  $16 \times 8$  data acquired 2 h after the  $8 \times 8$  data.

**Table 1**  
Comparison of SNR and metabolite peak ratios for normal vs. compressed sensing phantom data

	Peak SNR	Ala/Lac ratio	Pyr/Lac ratio	Pyr-H <sub>2</sub> O/lac ratio
Trial 1				
Normal $16 \times 8$	63.2	.55	.46	.43
Compressed sensing $16 \times 8$	64.8	.55	.47	.41
Trial 2				
Normal $16 \times 8$	59.7	.94	.50	.47
Compressed sensing $16 \times 8$	68.7	.89	.46	.42
Trial 3				
Normal $16 \times 8$	52.5	.66	.40	.40
Compressed sensing $16 \times 8$	62.4	.60	.37	.36

characteristic of cancer tissue in the TRAMP model. One difference is that due to lower starting SNR, residual coherent aliasing, and the sparsifying effect of the  $\ell_1$  reconstruction, the  $16 \times 8$  mouse data do not show tiny peaks such as the alanine and pyruvate-H<sub>2</sub>O bumps seen in the single spectrum of Fig. 6a. The reason is that high contrast spectral peaks result in large distinct sparse coefficients which can be recovered even when vastly undersampled, whereas very small peaks close to the noise floor could be submerged by both apparent noise caused by aliasing [10] and the true underlying noise that they would not be recoverable. As discussed in the artifacts section of [10], with increased undersampling the most distinct artifacts in compressed sensing are not the usual loss of resolution or increase in aliasing, but the loss of very small peaks. For the specific data shown in Fig. 6, the SNR of the alanine and pyruvate-H<sub>2</sub>O bumps in the  $8 \times 8$  acquisition were 17.7 and 14.3. By going to half the voxel size with  $16 \times 8$  resolution, the true SNR in each voxel would be halved, in other

words reduced to about 8. In addition, the extra apparent noise caused by undersampling and reconstruction inaccuracies [10] would further hurt the  $\ell_1$  reconstruction. According to the literature, coefficients can be recovered up to a multiple of the noise variance [29,30], and Lustig et al., describes this phenomenon for MRI in more detail while providing some examples [10]. For our spectra, we noticed that we needed an SNR of about 6–7 to distinguish a spectral peak from random spikes in the noise floor. Therefore, we believe the disappearance of the small peaks can be attributed to their being too close to the noise floor. In this scenario, no amount of denoising will recover the peaks, which is an important limitation of compressed sensing. Compressed sensing works best for applications with high SNR that are acquisition time limited such as measuring the lactate/pyruvate ratio in tumors. The second and third in vivo trials showed similar results except for a lower lactate/pyruvate ratio because of the less advanced disease stage of those tumors. Table 2 gives a quantitative comparison of the

**Table 2**  
Comparison of SNR and metabolite peak ratios for normal vs. CS mouse data

	Peak SNR	Lac/Pyr ratio	Standard deviation of Lac/Pyr ratio
Trial 1			
Normal $8 \times 8$	133.9	2.44	.432 ( $n = 16$ voxels)
Compressed sensing $16 \times 8$	107.4	2.51	.558 ( $n = 14$ voxels)
Trial 2			
Normal $8 \times 8$	41.1	1.54	.315 ( $n = 4$ voxels)
Compressed sensing $16 \times 8$	34.3	1.29	.353 ( $n = 8$ voxels)
Trial 3			
Normal $8 \times 8$	75.3	1.36	.246 ( $n = 4$ voxels)
Compressed sensing $16 \times 8$	64.8	1.38	.522 ( $n = 8$ voxels)

two reconstructions for each trial, showing SNR and lactate/pyruvate peak ratios. Halving the voxel size would normally reduce SNR by a factor of 2, but due to  $\ell_1$  denoising and apodization, as discussed in the previous section, the final SNR for the  $16 \times 8$  data was only 14–20% lower than that of the  $8 \times 8$ . Finally, as shown in Table 2, the ratios and standard deviations of the ratios match up reasonably well.

#### 4.3. Limitations and future work

The parameters controlling the level of denoising, and thus the final SNR, in this work were chosen manually. The parameters were chosen to strike a balance between denoising and the data fidelity constraint of Eq. (2). In the future, an automatic parameter choice scheme would be desirable. In addition, by employing more of the techniques in [10], such as phase estimation and variable density sampling patterns, it should be possible to perform more denoising and recover a higher SNR percentage without sacrificing data fidelity. Since the distribution of energy in  $k$ -space is localized close to the  $k$ -space origin, variable density sampling corresponding to that distribution has a better initial signal to aliasing interference ratio than uniform undersampling. It was reported in [10] that variable density schemes result in faster convergence and significantly better overall reconstruction quality and we expect our case to perform in the same way. In addition, investigation into different wavelet sparsifying transforms could yield further performance enhancements. Building on the blipped methodology presented in this paper to develop a sampling pattern in which  $k_f$ ,  $k_x$ , and  $k_y$  are undersampled could also provide a substantial performance gain by exploiting 3D sparsity and spreading aliasing into three dimensions, thus producing more incoherent aliasing. Lastly, combining parallel imaging and compressed sensing could be a viable avenue of investigation. With the abovementioned developments, much higher rates of acceleration might be obtainable in future hyperpolarized C-13 metabolic imaging studies.

## 5. Conclusions

An initial design and results for hyperpolarized  $^{13}\text{C}$  compressed sensing were presented. Key to the design was the exploitation of sparsity in hyperpolarized spectra and an implementation that used blips to undersample in  $k_f$  and  $k_x$ . Phantom experiments showed low SNR loss while preserving accuracy of metabolite peak ratios, and mouse trials demonstrated the in vivo feasibility of improving spatial resolution without increasing scan time in hyperpolarized  $^{13}\text{C}$  flyback 3D-MRSI. In addition, we discussed the unique properties of the  $\ell_1$  reconstruction in compressed sensing, such as wavelet denoising and the tendency to lose low contrast features such as small peaks. This study has demonstrated feasibility and potential value for applying compressed sensing to hyperpolarized  $^{13}\text{C}$  spectroscopy, and with further technique development even better performance is expected.

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